## Regularization of an autoconvolution problem in ultrashort laser pulse characterization

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JOHANNES KEPLER UNIVERSITY LINZ
$\square$ Introduction
$\square$ SD-SPIDER method
$\square$ Mathematical Analysis
$\square$ Discretization
$\square$ Regularization
$\square$ Numerical results

## Overview

$\square$ Introduction

SD-SPIDER method<br>Mathematical Analysis<br>Discretization<br>Regularization<br>Numerical results

## Motivation

Why study ultra-short laser pulses?
to create shorter, stronger pulses; to enhance optical systems; medicine, material processing, etc.

Problem: measurements limited by electronics (order $10^{-12} \mathrm{~s}$ ) Development of pulse durations:


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Solution: sample pulse by itself

## Laser pulse representation

Time domain: electric field $E(t)$, envelope $A(t)$, intensity $I(t)=|A(t)|^{2}$


Fourier domain: amplitude $\mathcal{A}(\omega)$, phase $\varphi(\omega)$, spectrum $\mathcal{I}(\omega)=|\mathcal{A}(\omega)|^{2}$

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$\square$ SD-SPIDER method

Mathematical Analysis

Discretization

Regularization

Numerical results

- SD-SPIDER=

Self-Defraction Spectral Phase Interferometry for Direct Electric-field Reconstruction

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- theory presented at "Conference on Lasers and Electro-Optics", 2010
- reasons for introduction: applicable for ultraviolet radiation, good signal strength because it uses third-order optical effects


## basics of nonlinear optics

- Polarization $\tilde{P}$ caused by an electric field $\tilde{E}$,

$$
\tilde{P}(t)=\epsilon_{0}\left[\chi^{(1)} \tilde{E}(t)+\chi^{(2)} \tilde{E}^{2}(t)+\chi^{(3)} \tilde{E}^{3}(t)+\ldots\right]
$$

may act as source of electromagnetic radiation:

$$
\nabla \times(\nabla \times E)+\frac{n^{2}}{c^{2}} \partial_{t}^{2} E=-\mu_{0} \partial_{t}^{2} P_{\mathrm{NL}}(E)
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■ Refraction index $n$ and Kerr-effect:

$$
n(\omega)=n_{0}+n_{2}|E(\omega)|^{2}
$$

(each frequency is refracted slightly differently)
$\chi^{(3)}$-media allow a four-wave mixing process


## Principle


k-vector-diagram:

$\overrightarrow{\Delta k}\left(\omega_{S D}, \omega_{\mathrm{p}}, \omega_{\mathrm{cw}}\right)$

$$
=-\vec{k}_{c w}\left(\omega_{c w}\right)+\vec{k}_{p}\left(\omega_{\mathrm{p}}\right)+\vec{k}_{p}\left(\omega_{S D}+\omega_{c w}-\omega_{\mathrm{p}}\right)-\vec{k}_{S D}\left(\omega_{S D}, \omega_{c w}, \omega_{\mathrm{p}}\right) .
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$$

energy conservation $\omega_{\mathrm{p}}+\omega_{\mathrm{p}}=\omega_{S D}+\omega_{c w}$ still holds

## The autoconvolution effect

- pulses considered as plane waves:



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- a wave $p_{1}$ of each frequency creates an interference pattern with cw-wave
■ at each pattern, photons $p_{2}$ of each frequency are refracted
- SD-signal is sum of all combinations

$$
\mathcal{E}_{p}\left(\omega_{\mathrm{p}}\right) \mathcal{E}_{p}\left(\omega_{S D}+\omega_{c w}-\omega_{\mathrm{p}}\right) \mathcal{E}_{c w}
$$

## Equation in physical formulation

$$
\begin{aligned}
& \mathcal{E}_{S D}\left(\omega_{S D}\right)=\int_{0}^{\omega_{S D}+\omega_{c w}} \mathcal{K}\left(\omega_{S D}, \omega_{\mathrm{p}}\right) \mathcal{E}_{p}\left(\omega_{\mathrm{p}}\right) \mathcal{E}_{p}\left(\omega_{S D}+\omega_{c w}-\omega_{\mathrm{p}}\right) d \omega_{\mathrm{p}} \\
& \operatorname{supp} \mathcal{E}_{p}=\left[\omega_{\mathrm{p}}^{l}, \omega_{\mathrm{p}}^{u}\right], \operatorname{supp} \mathcal{E}_{S D}=\left[2 \omega_{\mathrm{p}}^{l}-\omega_{c w}, 2 \omega_{\mathrm{p}}^{u}-\omega_{c w}\right]
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$$
\begin{array}{r}
\mathcal{K}\left(\omega_{S D}, \omega_{\mathrm{p}}\right)=\frac{\mu_{0} c L}{2} \frac{\omega_{S D}}{n\left(\omega_{S D}\right)}
\end{array} \chi^{(3)}\left(\omega_{S D},-\omega_{c w}, \omega_{\mathrm{p}}, \omega_{S D}+\omega_{c w}-\omega_{\mathrm{p}}\right) .
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$\mathcal{K}$ continuous, complex valued

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$\mathcal{K}$ continuous, complex valued
unknown, so far neglected

## mathematical formulation

- after transformation and renaming:

$$
\begin{aligned}
y(s) & =F[x](s)=\int_{0}^{s} k(s, t) x(t) x(s-t) d t \\
y=F(x) \quad & 0 \leq t \leq 1,0 \leq s \leq 2
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available, possibly available, unknown

- $\varphi(t)=\varphi_{0}+\int_{-\infty}^{t} G D(\tau) d \tau$


## Does $B(s)$ provide important information?



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Yes, it does! Thus also $B(s)$ available as measurement.

■ measurements (indicated by.$\delta$ ) "close" to correct data, but not exact

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- measurements (indicated by $\cdot{ }^{\delta}$ ) "close" to correct data, but not exact
- $A^{\delta} \rightarrow A, B^{\delta} \rightarrow B, \psi^{\delta} \rightarrow \psi$ as $\delta \rightarrow 0$
- no information about size of error $\delta$ available
- Statement of the problem: given $A^{\delta}, B^{\delta}, \psi^{\delta}$ and $k(s, t)$, find $\varphi$ such that

$$
B^{\delta}(s) e^{i \psi^{\delta}(s)}=\int_{0}^{s} k(s, t) A^{\delta}(t) e^{i \varphi(t)} A^{\delta}(s-t) e^{i \varphi(s-t)} d t
$$

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## III-posedness

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F x=y, \quad F: L^{2}[0,1] \mapsto L^{2}[0,2]
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An operator $F$ is called ill-posed, if it violates at least one of

## Hadamard's conditions:

(a) for each given data $y$ there exists a solution $x$
(b) this solution is unique
(c) the solution depends continuously on the data

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## Hadamard's conditions:

(a) for each given data $y$ there exists a solution $x$
(b) this solution is unique
(c) the solution depends continuously on the data
(a) violated because $F(x) \in C_{\mathbb{C}}[0,2] \forall x \in L_{\mathbb{C}}^{2}[0,1]$

## Injectivity

■ for $k(s, t) \equiv 1$ and $k(s, t)=k(s): F\left(x_{1}\right)=F\left(x_{2}\right)$ has two solutions $x_{1}=x_{2}$ and $x_{1}=-x_{2}$ by Titchmarsh's theorem

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■ because of periodicity, $\varphi \equiv \varphi+2 \pi$

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■ nonlinear operator requires local analysis

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## Definition

We define an operator $\mathcal{F}, \mathcal{F}: \mathcal{X} \rightarrow \mathcal{Y}$ to be locally ill-posed in $x_{0} \in \mathcal{X}$ if, for arbitrarily small $\rho>0$ there exists a sequence $\left\{x_{n}\right\} \subset B_{\rho}\left(x_{0}\right) \subset X$ satisfying the condition

$$
\mathcal{F}\left(x_{n}\right) \rightarrow \mathcal{F}\left(x_{0}\right) \text { in } \mathcal{Y} \text { as } n \rightarrow \infty \text {, but } x_{n} \nrightarrow x_{0} \text { in } \mathcal{X} .
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## Theorem (Gorenflo \& Hofmann '94, adapted in Gerth '11)

The autoconvolution operator $F$ is everywhere locally ill-posed.
$\Rightarrow(c)$ is violated too! Regularization is necessary.

## Fréchet-derivative

The Fréchet-derivative of $F$ in a point $x_{0}$ is given by

$$
\left[F^{\prime}\left(x_{0}\right) h\right](s)=\int_{0}^{s}(k(s, t)+k(s, s-t)) x_{0}(s-t) h(t) d t
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although $F$ is in general non-compact, $F^{\prime}$ is always compact!

## Overview

Discretization$\square$ Regularization
Numerical results

■ equation: $y(s)=\int_{0}^{s} k(s, t) x(s-t) x(t) d t$
$■ \operatorname{supp} x=\left[t_{l}, t_{u}\right], \operatorname{supp} y=\left[2 t_{l}-t_{c w}, 2 t_{u}-t_{c w}\right]$

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$■ \operatorname{supp} x=\left[t_{l}, t_{u}\right], \operatorname{supp} y=\left[2 t_{l}-t_{c w}, 2 t_{u}-t_{c w}\right]$
■ discretization using rectangular rule

$$
y\left(s_{m}\right)=\sum_{j=1}^{N} k\left(s_{m}, t_{j}\right) x\left(s_{m}+t_{c w}-t_{j}\right) x\left(t_{j}\right) \Delta t
$$

with $\Delta t=\frac{t_{u}-t_{l}}{N-1}, t_{j}=t_{l}+(j-1) \Delta t, s_{m}=2 t_{j}+(m-1) \Delta t$
$y_{m}:=y\left(s_{m}\right), x_{n}:=x\left(t_{n}\right), k_{m, n}:=k\left(s_{m}, t_{n}\right)$
in matrix-form $\underline{y}=\underline{F}(\underline{x}) \underline{x}$, with
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Decomposition, with $\circ$ as element-by-element multiplication:
$\underline{F}=\underline{K} \circ \underline{X}$
analogously: Fréchet-derivative

$$
\left[\underline{F^{\prime}}\left(\underline{x_{0}}\right) \underline{h}\right]_{m}=\sum_{j=1}^{N}\left(k\left(s_{m}, t_{j}\right)+k\left(s_{m}, s_{m}+t_{c w}-t_{j}\right)\right) x_{0}\left(s_{m}+t_{c w}-t_{j}\right) h\left(t_{j}\right) \Delta t
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resulting matrix $\underline{F^{\prime}\left(x_{0}\right)}=\left(\underline{K}+\underline{K^{\prime}}\right) \circ \underline{X_{0}}$
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resulting matrix $\underline{F^{\prime}\left(x_{0}\right)}=\left(\underline{K}+\underline{K^{\prime}}\right) \circ \underline{X_{0}}$
advantage: time-consuming calculation of the matrices $\underline{K}$ and $\underline{K}^{\prime}$ has to be performed only once for each measurement setup

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$\square$ Mathematical Analysis

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## A Levenberg-Marquardt-Type approach

- we let the complete pulse $x$ be unknown, whereas $y$ is given
- Iteration rule:

$$
\underline{x}_{(l+1)}^{\delta}:=\underline{x}_{(l)}^{\delta}+\gamma\left(\underline{F}^{\prime}\left(\underline{x}_{(l)}^{\delta}\right)^{*} \underline{F}^{\prime}\left(\underline{x}_{(l)}^{\delta}\right)+\alpha \underline{L}^{*} \underline{L}\right)^{-1} \underline{F}^{\prime}\left(\underline{x}_{(l)}^{\delta}\right)^{*}\left(\underline{y}^{\delta}-\underline{F}\left(\underline{x}_{(l)}^{\delta}\right)\right.
$$

for $l=0, \ldots, l^{*}$, aimed at minimizing

$$
\left\|\underline{y}^{\delta}-\underline{F}\left(\underline{x}_{(l)}\right)-\underline{F}^{\prime}\left(\underline{x}_{(l)}\right)\left(\underline{x}-\underline{x}_{(l)}\right)\right\|^{2}+\alpha\left\|\underline{L}\left(\underline{x}-\underline{x}_{(l)}\right)\right\|^{2}
$$

$\underline{L}(\underline{x})$ approximating the second derivative of $x$

## A Levenberg-Marquardt-Type approach

- we let the complete pulse $x$ be unknown, whereas $y$ is given
- Iteration rule:

$$
\underline{x}_{(l+1)}^{\delta}:=\underline{x}_{(l)}^{\delta}+\gamma\left(\underline{F}^{\prime}\left(\underline{x}_{(l)}^{\delta}\right)^{*} \underline{F}^{\prime}\left(\underline{x}_{(l)}^{\delta}\right)+\alpha \underline{L}^{*} \underline{L}\right)^{-1} \underline{F}^{\prime}\left(\underline{x}_{(l)}^{\delta}\right)^{*}\left(\underline{y}^{\delta}-\underline{F}\left(\underline{x}_{(l)}^{\delta}\right)\right.
$$

for $l=0, \ldots, l^{*}$, aimed at minimizing

$$
\left\|\underline{y}^{\delta}-\underline{F}\left(\underline{x}_{(l)}\right)-\underline{F}^{\prime}\left(\underline{x}_{(l)}\right)\left(\underline{x}-\underline{x}_{(l)}\right)\right\|^{2}+\alpha\left\|\underline{L}\left(\underline{x}-\underline{x}_{(l)}\right)\right\|^{2}
$$

$\underline{L}(\underline{x})$ approximating the second derivative of $x$

- Questions:

■ how to choose $\underline{x}_{0}$ ?
■ how to choose $l^{*}$ ?
■ how to choose $\alpha$ ?

## Choice of $\underline{x}_{0}=A_{0} e^{i \varphi_{0}}$

obviously, $A_{0}:=A^{\delta}$
first idea for phase: $\varphi_{0}(t) \equiv 0$


$(\delta=0, \alpha=0)$
idea: calculate good guess. Observe

$$
\begin{aligned}
& B^{\delta}(s) e^{i \psi^{\delta}(s)}=\int_{0}^{s}|k(s, t)| A^{\delta}(t) A^{\delta}(s-t) e^{i\left(\varphi(t)+\varphi(s-t)+\phi_{\text {kernel }}\right)} d t \\
\Rightarrow & \text { set } \varphi_{0}(t)=\frac{1}{2}\left(\mathcal{P}_{s \mapsto t}(\psi(s))\right)-\phi_{\text {kernel }}\left(s^{*}, t\right) \text { for } s^{*} \text { fixed }
\end{aligned}
$$



problem for slightly changed fundamental phase


## best result with kernel correction




## best result with kernel correction



$\Rightarrow$ set starting phase to constant zero

## When to stop the iteration?

An example iteration:

| $(l)$ | $\\| \underline{\underline{E}\left(x_{(l)}^{\delta}\right)-\underline{y}^{\delta} \\|}$ | $\left\\|\left\|\underline{x}_{(l)}^{\delta}\right\|-A^{\delta}\right\\|$ |
| :---: | :---: | :---: |
| 1 | $9.5819 \mathrm{e}-01$ | 0.5252 |
| 20 | $2.4115 \mathrm{e}-02$ | 0.7916 |
| 40 | $2.0682 \mathrm{e}-02$ | 0.7937 |
| 60 | $1.5369 \mathrm{e}-02$ | 0.6077 |
| 100 | $1.3792 \mathrm{e}-03$ | 0.1964 |
| 120 | $1.1022 \mathrm{e}-03$ | 0.1701 |
| 140 | $9.4595 \mathrm{e}-04$ | 0.1623 |
| 143 | $9.2340 \mathrm{e}-04$ | 0.1622 |
| 144 | $9.1606 \mathrm{e}-04$ | 0.1623 |
| 150 | $8.7480 \mathrm{e}-04$ | 0.1632 |
| 250 | $3.1613 \mathrm{e}-04$ | 0.2020 |

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$\Rightarrow$ choose $l^{*}$ such that $\left|\left|\underline{x}_{(l)}^{\delta}\right|-\underline{A}^{\delta} \|\right.$ is minimal

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- instead, make use of $A^{\delta}$ again: choose $\alpha^{*}$ such that

$$
\left|\left|\left|\underline{x}_{\alpha^{*}}^{\delta}\right|-A^{\delta} \|=\min _{n}\right|\right|\left|\underline{x}_{\alpha_{n}}^{\delta}\right|-A^{\delta}| |
$$

## Overview

Numerical results

## A very smooth fundamental pulse




## SD-pulse, 5\% relative noise added




## reconstruction, $\alpha=5.86 \cdot 10^{6}$




## A more oscillating pulse




## noise-free SD-pulse




## reconstruction, $\alpha=2.17$




## reconstruction, $1 \%$ relative noise in data




## Real data situation

unfortunately, no results available. Main reasons:

- measurements without magnitudes
- unknown factor in model $\Rightarrow$ error in the model
- frequency domains of $x$ and $y$ do not match
D. Gerth, B. Hofmann, S. Birkholz, S. Koke, and
G. Steinmeyer Regularization of an autoconvolution problem in ultrashort laser pulse characterization, submitted
D. Gerth, Regularization of an autoconvolution problem occurring in measurements of ultra-short laser pulses, Diploma thesis, Chemnitz University of Technology, Chemnitz, 2011, http://nbn-resolving.de/urn:nbn:de:bsz:ch1-qucosa-85485.
- R. Gorenflo, B. Hofmann, On autoconvolution and regularization, Inverse Problems 10 (1994), pp. 353-373.

囯 S. Koke, S. Birkholz, J. Bethge, C. Grebing, G. Steinmeyer, Self-diffraction SPIDER, Conference on Laser and Electro Optics (CLEO), San Jose, CA, 2008.
D. Gerth, B. Hofmann, S. Birkholz, S. Koke, and
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D. Gerth, Regularization of an autoconvolution problem occurring in measurements of ultra-short laser pulses, Diploma thesis, Chemnitz University of Technology, Chemnitz, 2011, http://nbn-resolving.de/urn:nbn:de:bsz:ch1-qucosa-85485.

- R. Gorenflo, B. Hofmann, On autoconvolution and regularization, Inverse Problems 10 (1994), pp. 353-373.
(i) S. Koke, S. Birkholz, J. Bethge, C. Grebing, G. Steinmeyer, Self-diffraction SPIDER, Conference on Laser and Electro Optics (CLEO), San Jose, CA, 2008.

Thank you for your attention! Are there any questions?

